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# The X-ray Crystal Structure of Diethylammonium Tetracyanopalladate

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The crystal structure of  $[(C_2H_5)_2NH_2]_2Pd(CN)_4$  has been determined by inspection of Patterson bidimensional synthesis (*hk*0) and (*h*0*l*) and by refinement of tridimensional intensity data. The space group of the primitive cell is  $P\overline{1}$ , but a more convenient double cell has been chosen whose space group is  $C\overline{1}$ , with  $a=15\cdot79\pm0\cdot01$ ,  $b=9\cdot10\pm0\cdot01$ ,  $c=6\cdot35\pm0\cdot01$  Å,  $\alpha=84^{\circ}11'\pm1'$ ,  $\beta=92^{\circ}42'\pm1'$ ,  $\gamma=94^{\circ}5'\pm1'$ . The complex anion is centred on the origin and the distance between two heavy atoms is considerably larger than the corresponding one in the inorganic salts of the  $Pd(CN)_4^2$  anion. The latter has the well known square-planar configuration. The four CN ligands are bound to the adjacent cations by H bonds of the N-H···N type, so forming an infinite chain lying in the *b* axis direction. Neighbouring chains in the ( $\overline{2}01$ ) direction form planes whose inclination approximates that of the complex group. The stacking of the planes performs the tridimensional structure. The weak van der Waals bonds between chains and between planes is responsible for the important thermal motion of the atoms.

#### Introduction

The study of the inorganic salts of  $Pd(CN)_4^2$  and  $Pt(CN)_4^2$  complexes in the solid state started with the structural works of Brasseur & de Rassenfosse (1935, 1936, 1937, 1938). Recently, the tridimensional X-ray analysis of CaPd(CN)\_4.5H\_2O (Fontaine, 1968) and Rb\_2Pd(CN)\_4.H\_2O (Dupont, 1970) confirmed some structural features that are believed characteristic of the complex group present in the crystal lattice.

In the crystals just mentioned, the square-planar ions are always found in columns parallel to the *c*-axis direction. The shortest distance between two neighbouring heavy atoms varies from  $3 \cdot 13$  Å in the heptahydrate of Mg<sup>2+</sup> to  $3 \cdot 72$  Å in the monohydrate of Rb<sup>+</sup>. Optical properties such as fluorescence emission (Pt compounds only), ultraviolet and visible absorption seem to be related to the heavy-metal spacing (Fontaine, Moreau & Simon, 1968). In order to increase the latter, we have replaced the metallic cation by a large organic one and we have studied the structure and properties of such a compound.

Previous preliminary research (Jérôme-Lerutte, 1967) allows us to choose the diethylammonium salt as a specially demonstrative example.

#### **Crystal data**

The cell dimensions and space group of the diethylammonium tetracyanoplatinate have already been published (Jérôme-Lerutte, 1967). The palladate is quite isotypic. It is triclinic;  $a=15.79\pm0.01$ , b=

H F2 FC

K =

13579113571-35791135

K3 -2 -4 -10 -12 -14 -16

K\* C?448C2462462 11112462468 -112462462 
 kt
 --, (2)
 3

 0
 0.01
 6.37

 2
 7.64
 2.90

 4
 7.63
 2.90

 4
 7.63
 2.90

 4
 7.63
 2.90

 4
 7.63
 2.90

 4
 7.63
 2.90

 4
 7.73
 3.60

 10
 16.4
 16.4

 -10
 16.4
 16.4

 -11
 16.4
 16.7

 -12
 14.6
 16.7

 -6
 3.61
 3.64

 -10
 16.4
 16.7

 -12
 14.6
 16.7

 -14
 .0
 113

 Kz
 .7, 1.4
 .7

 $9.10 \pm 0.01$ ,  $c = 6.35 \pm 0.01$  Å,  $\alpha = 84^{\circ} 11' \pm 1'$ ,  $\beta =$  $92^{\circ}42' \pm 1'$ ,  $\gamma = 94^{\circ}5' \pm 1'$ ,  $D = 1.305 \pm 0.010$  g.cm<sup>-3</sup>, Z=2. The cell just described is doubly primitive with space group  $C\overline{1}$ . It was chosen because it makes the graphical representation of the Patterson and Fourier syntheses easier; also it permits interesting comparisons to be made with the monoclinic cell of the triethylammonium compound whose structure will be published later. No physical test has been performed regarding the symmetry center but the good results obtained with that assumption are sufficient to prove its presence.

# Table 1. Observed and calculated structure factors

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•	121	715	٠.	· , · •		2	2. 1 22.	-9	434 4	677 -1 646	5 1#3	179	-12 379 -14 190	344
4	575	277 427	~	\$16		7	2 20 1211. 6 24 - 376	-11	414 319	303 397 K	-7, L:	1	-16 187 -18 132	1 72
12	677	47	2	64.		11	642 657 452 372	-15	95	154	1 427	405	K= 1, 1=	z
14	716	24.5	۴ ۹	297	377	15	255 261 195 181	<	4. (=	,	5 347	296 251	1 439	392
20	1.4	177	10	215	177	1.1	7 1 481		1	970 I	0 286	25R	3 751	737
Ke	1. 1.		14	105	213		512 442	2	465	44 1	3 262	241	7 482	450
	619		12	191	100	-11	415 389	10	3.96	374 -	1 454	439	11 511	462
ì	90	1156	-6	225	419	-14	214 214	- 2	879		5 349	344	15 226	231
,	204	254	-10	240	710	-12	149 155		207	310 -1	9 185	191	-3 991	1007
11	344	367	-14	25	245	К	1+ L+ - 1	-10	500	480 -1	3 159	155	-9 405	399
15	327	2 15	-15	2.0	10	1 1	210 1024	-14	270 2	62 ×=	8, L.	1	-13 251	242
19	122	123	K.=	7, 19	Ş	3	503 594	-15	94 1	134	0 389	352	-15 213	200
-5	703	P19	3	3 **	6.00		393 389		4, L.	1 -	4 320	295	-19 81	103
-9	5R4	543	7	220	216	13	372 373	0	389 4	- 101	5 244 8 274	243	K1, L.	2
-13	344	378	11	232	231	17	293 279	4	542	552 -1 538 -1	2 229	230	1 623	689 761
-17	198	274	15	120	150	-1 1	216 1193	â	573	708 -1 581	4 172	172	7 456	423
-10	109	133	-1	501 439	497	-3	59C 723	10	135	398 X≖ 14∩	-8, L=	1	9 649 11 460	444
K=	2, 14	0	-7	2 96	279	-7	525 524 339 381	14	265 238 2	244	0 283 2 301	276 271	13 314 15 242	320 252
2	352	1124 366	-11	371 228	281 271	-11	256 323 326 308	18	173 1	529	4 263 6 379	245 326	17 187 19 92	198 124
6	945 866	912 821	-13 -15	261 201	249	-17	190 205 158 167	-6 -8	320 384 3	317 387 1	8 353 0 242	312 217	-1 743 -1 789	816 773
10	583 354	554 324	K=	8, L=	0	-19	129 144	-1C -12	231 2	639 1 237 1	2 146	145 189	-5 292 -7 586	294 610
12	376 275	358 247	с	392	417	K =	2, 1. 1	-14	216 2	218 - 134 -	2 385	364 316	-9 567 -11 258	503 256
16 18	209	195	2	340 229	344 256	21	272 247	χ	5. 1.	, :	6 201 8 177	205 178	-13 187 -15 176	200
-2 -4	219	206	6	250	250	6	883 944 576 548	ı	645 6	-1	2 221	208	-17 141	158
- 6 - 8	803	995 504	10 12	199	1 42	я 10	265 231	3	481 4	67 59 K	9. L.	1	K- 2.L-	2
-10	314	314	-2	294	299	12	345 350 275 266	-3	524 551 5	534	1 236	217	0 1123 1	-
-14	272	276	-6	289	263	15	197 199	-7	450 4	25 -	3 284	246	2 850	834
-18	162	157	-10	274	253	-4	549 503	-11	353 3	351 -	7 246	225	6 396	413
×	3, L=	0	-14	100	137	-8	698 717	-15	165 1	67 -L	1 232	217	10 367	358
1	342 731	322	K =	9, i=	0	-12	219 222	** -	5. 13	1 1 1	-9.14		14 234	229
5	856	837	1	244	231	-16	190 194	· · ·	530 6	1		747	-2 535	534
9	618	571	5	253	253		2. 1. 1	ż	374 3	37R	3 305	275	-6 502	506
13	213	202	11	203	203		336 312	7	498 4	89	7 278	246	-10 386	371
17	162	154	-1	278	269	ě	360 333	11	371	346 1	1 190	168	-14 167	158
-1	394	373	-5	25A	233	6	461 436	15	216 2	208 -	3 187	179	-18 133	137
- 5	808	865	-9	295	268	12	610 588	-1	573	- 562	7 237	219	K= -2, L=	2
-9	615	602	-13	104	1 32	14	311 294		461 4	178 -1	104	137	0 973	941
-13	276	277	Ke	10, L=	0	14	159 168	-9	259 2	260 K=	10, L=	1	2 696	730
-17	154	152	0	148	155	-4	491 523	-13	228 2					473
-17		124	-	210	210	-0	0.1 102	-16	164		4 172	162	4 552 6 503	403
<u>.</u>				234		-8	656 659	-15	156 1		4 172 6 253 8 272	162 226 238	4 552 6 503 8 567 10 362	602 349
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$\begin{array}{c} 0 & 2.77 \\ 2.72 \\ 3 & 765 \\ 1 & 3 \\ -5 \\ 1 & -5 \\ $	286	-9, L= -9, L= -9, L= 267 135 135 135 135 135 135 14 137 14 107 127 17, L= 127 17, L= 127 17, L= 127	21967 11654	4 6 8 16 2 14 1 - 2 4 1 1 14 - 2 4 1 1 14 - 2 4 1 14 - 2 4 1 14 - 2 1 14 - 1 1 - 1 15 - 7 - 0 1 - 1 15 - 7 - 0 1 - 1 15 - 0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 - 0 -	431 544270956 3 3 7444972743223115 64447785469865 1 9460778544886623115	4431983974 7 44457.000 7 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	n2469024472 k 1367911136703	467 467 467 407 407 407 407 407 407 407 40	499933232134 441221641321 99933127001 269632932127001 2696329322144	10 12 12 12 12 12 12 12 12 12 12	299 189 317 317 24, Lx 314 323 282 314 323 282 315 46 147 147 7, L= 306	313           2194           309           313           3294           309           313           329           3329           3329           3329           3329           3329           3329           3329           3329           313           329
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# Experimental

Intensity data were recorded photographically with Cu  $K\alpha$  radiation by the multiple-film technique using an integrating Weissenberg camera. The intensities corresponding to the reciprocal lattice planes (hk0), (hk1), (hk2), (h0l), (h1l), (h2l) and (h3l) were investigated and measured with a Unicam microdensitometer. Eighty-seven unknown parameters, including 60 anisotropic thermal motion factors, were determined from the intensity data. The proportion of accidentally extinguished intensities was quite low because the

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7 711 9 195 -1 262 -3 234 -5 146 -9 154 -11 111 1131 730 3 249 5 259 7 203 7 203 9 211 9 204 -1 263 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 164 -1 111 -1 264 -1 111 -1 264 -1 111 -1 264 -1 111 -1 270 -1 264 -1 264 -1 111 -1 270 -1 264 -1 264 -1 111 -1 270 -1 264 -1 270 -1 264 -1 270 -1 264 -1 270

heavy metal being at the cell origin influenced all reflexions equally. They were not included in the refinement as they could not change the structure results. A few experimental data close to the reciprocal lattice origin were also discarded since they were affected by secondary extinctions. The intensities of the 1169 remaining reflexions were corrected for the Lorentz and polarization effects as well as for absorption. An IBM 7090 computer program written by Dideberg (1966) was used for the latter correction.

### Solution and refinement of the structure

A first attempt to locate C and N atoms was performed by means of (hk0) and (h0l) Patterson maps (Figs. 1) & 2). The heavy atom is necessarily situated on the symmetry centre and CN ligands appear in the peaks nearest the heavy atom. The diethylammonium ion is more difficult to locate but its location is obtained by means of stereochemical considerations. After a few trials, the Pd, C and N positional parameters and isotropic temperature factors were included in a refinement cycle. The corresponding R factor (10.9%)doubtless shows the accuracy of the atoms ordering in the unit cell. The whole set of measured intensities was refined by the full-matrix least-squares method with anisotropic thermal motion factors for the atoms including hydrogen atoms. After a few cycles, an ultimate R factor of 5.5% was obtained and convergence was reached for Pd, C and N parameters. The leastsquare calculation was performed on an IBM 7040 computer using program ORFLS (Busing, Martin & Levy, 1962). The observed and calculated structure factors are listed in Table 1. Atomic diffusion factors are taken from tables published by Cromer & Waber (1965). Coordinates and temperature factors of nonhydrogen atoms are described in Table 2.

# Determination of the hydrogen atomic positions

In the first step of the work, H-atom contributions were neglected but the estimated C-C and C-N distances were systematically longer than the average literature values. H atoms were then introduced in hypothetical positions according to the following principles: C and

N<sup>+</sup> atom hybridization is tetrahedral, C-H distance is 1.09 Å, N-H(1) is 1.01 Å. Thus 6 of the 12 H atoms



2 sin β

Fig. 1. Patterson map (001).



Fig. 2. Patterson map (010).

Table 2. Coordinates and anisotropic temperature coefficients of non-hydrogen atoms (e.s.d.'s in parentheses)

		<i>T</i> =	$= \exp\left[-(\beta_{11}h^2 +$	$+\beta_{22}k^2+\beta_{33}l^2$	$+\beta_{12}hk+\beta_{13}$	$hl+\beta_{23}kl)]$			
	$\frac{x}{a}(104 \sigma x)$	$\frac{y}{b}(10^4 \sigma y)$	$\frac{z}{c}(10^4 \sigma z)$	$10^4 \beta_{11}(\sigma_{11})$	$10^4 \beta_{22}(\sigma_{22})$	10 <sup>4</sup> β <sub>33</sub> (σ <sub>33</sub> )	$10^4 \beta_{12}(\sigma_{12})$	$10^4 \beta_{13}(\sigma_{13})$	$10^4 \beta_{23}(\sigma_{23})$
Pđ	0	0	0	30.3 (0.3)	69.9 (1.2)	208 (2.9)	4.6 (0.4) -	-25.4 (0.6)	-17.1 (1.4)
·C(1)	0.0687 (7)	0.1513 (13)	0.1424 (18)	54 (4)	128 (14)	416 (32)	21 (7)	-46 (10)	28 (22)
C(2)	0.0656 (6)	0·8438 (12)	0.1641 (16)	50 (4)	105 (12)	349 (27)	-6(6)	-29 (8)	- 59 (20)
C(3)	0.1078 (17)	0.5519 (26)	0.7803 (26)	164 (21)	350 (44)	460 (50)	77 (25)	54 (25)	-20(43)
C(4)	0·1140 (12)	0.4333 (19)	0.6562 (24)	109 (10)	199 (25)	474 (45)	-3 (13)	62 (18)	-24 (34)
C(5)	0.2548 (8)	0.4982 (20)	0.4939 (27)	48 (5)	274 (30)	740 (74)	7 (10)	- 54 (16)	6 (42)
C(6)	0·2969 (11)	0.5207 (27)	0.2813 (35)	69 (8)	398 (50)	852 (88)	-9 (16)	79 (23)	-4 (54)
N(1)	0.1107 (8)	0.2432(13)	0.2202 (21)	82 (6)	150 (17)	687 (47)	2 (8)	-125 (14)	-87 (25)
N(2)	0.1030 (7)	0.7523 (11)	0.2574 (16)	72 (5)	143 (14)	464 (31)	35 (7)	- 59 (10)	35 (20)
N(3)	0.1619 (5)	0.4784 (10)	0.4597 (13)	51 (3)	135 (12)	343 (22)	16 (5)	-23 (7)	-4 (16)

were unambiguously located, the 6 others occupying vertexes of tetrahedra whose orientation is not known. Their positions were chosen by accounting for possible sterical hindrances. Later on, attempts were made to introduce the positional parameters of H atoms as variable quantities. The temperature factors were equal to the last isotropic factor of the C or N atom to which the H atom is attached. Positional errors were important and convergence was not reached for all atoms. But the hypothesis of an improvement of the proposed model was acceptable according to the Hamilton (1965) significance tests.

Hydrogen coordinates and their temperature factors are listed in Table 3.

## Table 3. Hydrogen coordinates and temperature-factor coefficients

	$T = \exp\left(-B \frac{\sin^2\theta}{\lambda^2}\right).$									
	$\frac{x}{a}$ (10 <sup>3</sup> $\sigma x$ )	$\frac{y}{b}$ (10 <sup>4</sup> $\sigma y$ )	$\frac{z}{c}$ (10 <sup>4</sup> $\sigma z$ )	Т						
H(1)	0.142 (9)	0.379 (16)	0.373 (19)	5.14						
H(2)	0.134(9)	0.577 (16)	0.421 (19)	5.14						
H(3)	0.146 (12)	0.320(22)	0.758 (27)	9.13						
H(4)	0.047 (12)	0.386 (20)	0.598 (25)	9.13						
H(5)	0.242(11)	0.630 (19)	0.626 (23)	7.31						
H(6)	0·272 (11)	0.399 (19)	0.643(23)	7.31						
H(7)	0·108 (14)	0.660 (25)	0.656 (31)	10.84						
H(8)	0.182(15)	0.610 (26)	0.884 (31)	10.84						
H(9)	0.109 (15)	0.508 (24)	0.875 (29)	10.84						
H(10)	0.356 (13)	0.582(24)	0.260 (29)	9.86						
H(11)	0.262 (13)	0.636 (23)	0.166 (28)	9.86						
H(12)	0.320 (13)	0.403(24)	0.168 (29)	9.86						



🚦 sin y

Fig. 3. Projection of the structure on the (001) plane.

#### Description of the structure

Bond lengths, bond angles and their associated standard deviations are listed in Tables 4 and 5. Projections of the structure on the (001) and (010) planes are shown in Figs. 3 and 4 respectively.

# Table 4. Bond length (e.s.d.'s in parentheses)

Å

Pd——C(1)	1.970 (15)
Pd - C(2)	2.001 (12)
C(1) - N(1)	1.173 (16)
C(2) - N(2)	1.156 (13)
C(3) - C(4)	1.419 (26)
C(5) - C(6)	1.530 (25)
C(4) - N(3)	1.504 (17)
C(5) - N(3)	1.478 (15)
C(3) - H(7)	1.2 (2)
C(3) - H(8)	1.4 (2)
C(3)–H(9)	0.7 (2)
C(4)–H(3)	1.2 (2)
C(4)–H(4)	1.2 (2)
C(5)–H(5)	1.6 (2)
C(5)–H(6)	1.3 (2)
C(6)–H(10)	1.1 (2)
C(6)–H(11)	1.4 (2)
C(6)–H(12)	1.4 (2)
N(3)-H(1)	1.13 (15)
N(3)–H(2)	1.03 (14)

Table 5. Bond angles (e.s.d.'s in parentheses)

0.4)
0·4́)
1.2)
0·9)
1.5)
1.1)
1.3)
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2)
2)
b) –
5)
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The complex group has the well known squareplanar configuration. Pd-C and C-N distances are in accordance with the values listed by Fontaine (1968) and Dupont (1970). The  $Pd(CN)_4^{2-1}$  ion is only slightly distorted from the theoretical  $D_{4h}$  symmetry. The mean plane calculated by the program of Pippy & Ahmed (1966) has the following equation:

$$0.686x - 0.015y - 0.727z = 0$$
.

The planarity is unquestionable since no out-of-plane distortion exceeds  $3\sigma$ . The group is approximately situated in the  $(\overline{2}01)$  plane. The diethylammonium cation lies approximately in a plane parallel to (010) and distant  $\frac{1}{2}b$  from the origin. C-N sp<sup>3</sup> distances agree with the average experimental value and with the theoretical one: 1.47 Å. The situation is not so clear for C-C distances: one value agrees fairly well with the theoretical 1.55 Å value but the other is wrong: the discrepancy is beyond the  $3\sigma$  limit. Such a contraction has no chemical explanation since the measured length should be about an  $sp^2$  C-C distance. Neither is the presence of systematic errors an explanation as no reflexion is specially affected by the C(3) atom displacement in its assumed position. A unique least-squares refinement brings it back in its place. Since thermal motion amplifies itself considerably at the cationic chain bottoms, the mean position could possibly be erroneous but only a low-temperature study should eliminate that kind of error.

# H bonds

Two H atoms are attached to the N(3) atom and point towards the N atoms of the CN ligands. The measured distances and angles are listed in Table 6.

Table 6. Distances and angles in the H-bond system

$N(1) \cdots H(1)$	1·69 (0·15) Å
$N(1) \cdots N(3)$	2.82 (0.02)
$N(2) \cdots H(2)$	1.90 (0.14)
$N(2) \cdots N(3)$	2.89 (0.02)
$N(1) \cdots N(3) \cdots N(2)$	110.5 (0.4)
$N(1) \cdots H(1) \cdots N(3)$	174 (10)
$N(2)\cdots H(2)\cdots N(3)$	159 (10)
$\mathbf{C}(1) \equiv \mathbf{N}(1) \cdots \mathbf{N}(3)$	161·8 (1·2)
$C(2) \equiv N(2) \cdots N(3)$	166-2 (0-9)

The N(1)...N(3) and N(2)...N(3) distances confirm the presence of two H bonds of the N-H...N type. According to Pimentel & McClellan (1960), the mean distance of such bonds is  $3 \cdot 10$  Å. Therefore, here they seem stronger than the average one. The accuracy of H positions is not sufficient to state the linearity of these bonds. In the large error limits, they do not seem highly bent. The ratio of H donor (NH bond) to H acceptor (CN ligand) is unity. So all N atoms in the



Fig. 4. Projection of the structure on the (010) plane.

structure take part in H-bond formation and contribute to the building of infinite chains lying in the b-axis direction (Fig. 5).

#### Packing

The chains parallel to the *b* axis lie side by side in the  $(\overline{2}01)$  plane with a displacement of  $\frac{1}{2}b$  between consecutive chains, as shown in Fig. 5. Fig. 6 shows the projection of the chains on the  $(\overline{2}01)$  plane, with their



Fig. 5. Schematic representation of the chain formation along the b axis direction in the ( $\overline{2}01$ ) plane.



Fig. 6. Projection of the chains on the  $(\overline{2}01)$  plane.

insistent sterical occupation. The outer circles in thick full line represent van der Waals spheres, whereas the inner circles in thin full line indicate the covalence spheres. Four contact distances between chains have been calculated, *i.e.* distances approximately equal to the sum of the van der Waals radii. Fig. 7 shows the stacking of ( $\overline{2}01$ ) planes through a projection of a (010) section between about x = -0.25 and x = 0.25. The complex groups are centered at the four corners of the cell. Between them, two cationic unities lie in the normal direction to the  $(\overline{2}01)$  plane. One ethyl group touches the  $Pd(CN)_4^2$  ion; two contact distances have been found at this place. Another plane-plane contact probably exists at the other end of the diethylammonium group between two H atoms. Table 7 lists the contact distances between non-bonded atoms.

# Table 7. Contact distances between non-bonded atoms (e.s.d.'s in parentheses)

	Symmetry	operators	
	x	У	Z
n <sup>I</sup> n <sup>II</sup>	$\frac{\frac{1}{2}-x}{\frac{1}{2}-x}$	$\frac{1}{2} - y$ $1 - y$	$\begin{array}{c}1-z\\1-z\end{array}$

	Table 7	1 (cont.)	
n <sup>III</sup> n <sup>IV</sup> n <sup>V</sup>	$ \begin{array}{c} x \\ \frac{1}{2} - x \\ - x \end{array} $	$\frac{y}{1-y}$	-1+z $-z$ $1-z$
	Contacts bet	tween chains	
C(1)-	H(6 <sup>I</sup> )	2.85 (0	)·17) Å
N(1)-	$H(6^{1})$	2·42 (0	)·17)
H(3)-	$H(12^{I})$	2.21 (0	)•26)
H(8)-	$H(11^{11})$	2.12 (0	)·29)
	Contacts bet	tween planes	
C(1)-	$H(12^{IV})$	2.92 (0	•20) Å
N(2)-	$-H(4^{v})$	2.75 (0	)•19)
H(11)	-H(8 <sup>111</sup> )	2.10 (0	)·28)

# Thermal motion

The principal axis of vibration ellipsoids is given in Table 8. The reference system of axis is defined for each particular case, except for all atoms of the complex group whose system is the same. In this latter group, the atoms preferably move in a direction perpendicular to the group plane. The amplitude of vibration in the group plane is nearly equal in all directions

Tuble 0, Therman encenteance for more of the order and	Table 8. Therma	l eigenval	ues for	non-hya	rogen	atoms
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 $\sqrt{O_i^2}$  are the r.m.s. amplitudes. Angles between the principal axis of the ellipsoids and reference axis are also given.

	i axis	$\sqrt{\bar{U}_{i}^{2}}$ (10 <sup>3</sup> )	1	il angle	m	im angle	n	in angle
Pd	1	0.162 (1)	Pd-C(1)	46 (9)°	PdC (1)	87 (1)°	⊥lm	136 (9)°
	2	0.170 (1)	,,	136 (9)	C(2') plane	88 (1)		134 (9)
	3	0.235 (1)	**	90 (1)	**	4 (1)	,,	86 (1)
<b>C</b> (1)	1	0.184 (13)	**	87 (6)	"	88 (4)	**	176 (6)
-(-)	2	0.261 (16)	**	166 (9)	**	104 (9)	**	94 (6)
	3	0.329 (12)	,,	104 (9)	**	14 (9)	**	89 (4)
C(2)	1	0.180 (18)	,,	7 (8)	,,	90 (5)	••	83 (8)
0(2)	2	0.245(12)	**	97 (8)	**	97 (11)	"	10 (9)
	3	0.293 (11)	**	91 (5)	**	7 (11)	••	83 (11)
N(1)	1	0.208(19)	,,	8 (10)	**	87 (3)	,,	97 (10)
1(1)	2	0.267(12)	"	82 (10)	**	95 (3)	••	9 (8)
	3	0.443 (13)	"	92 (3)	**	6 (3)	,,	85 (3)
N(2)	1	0.184(12)	,,	90 (4)	**	90 (3)	"	179 (4)
1.(2)	2	0.295(14)	**	177 (8)	**	93 (8)	**	90 (4)
	3	0.359 (11)	**	93 (8)	"	3 (8)	**	90 (3)
C(3)	1	0.297(17)	C(3) - C(4)	136 (18)°	$\perp C(3) - C(4) - N(3)$	81 (14)°	**	48 (14)
0(3)	2	0.346(22)	-(-),,-(-)	54 (18)	plane	48 (9)	**	63 (16)
	3	0.484 (28)	**	112 (7)	- >>	44 (8)	**	126 (6)
C(4)	1	0.281 (16)	C(4)-N(3)	32 (44)	$\perp C(3)-C(4)-N(3)$	97 (8)	,,	121 (44)
0(.)	2	0.297 (20)	· ,, ` ,	58 (44)	plane	85 (11)	**	33 (43)
	3	0.387 (18)	,,	86 (7)	,,	9 (9)	,,	98 (10)
N(3)	1	0.214 (9)	N(3)-C(5)	53 (9)	$\perp C(4) - N(3) - C(5)$	51 (10)	"	120 (6)
- (-)	2	0.253 (12)	· ,, · · ·	129 (11)	plane	41 (11)	••	80 (13)
	3	0·288 (́9) ́	**	60 (10)	**	78 (11)	**	33 (6)
C(5)	1	0.225 (13)	C(5)-N(3)	167 (6)	$\perp$ N(3)-C(5)-C(6)	83 (7)	,,	80 (4)
- (- )	2	0.334 (18)	***	80 (7)	plane	22 (12)	••	71 (12)
	3	0.411 (19)	••	83 (5)	**	110 (12)	,,	21 (11)
C(6)	1	0.266 (16)	C(6)C(5)	137 (5)	$\perp N(3)-C(5)-C(6)$	94 (6)	••	47 (5)
,	2	0·409 (24)	`,,``	77 (24)	plane	23 (35)	**	71 (25)
	3	0.439 (23)	"	49 (12)	**	113 (36)	,,	49 (16)

for the heavy atom, whereas the C atoms move with their minimum amplitude perpendicular to the Pd-C direction and the N atoms move along that direction. This inversion may be influenced by the presence of H bonds. In the cation, the motion is described with ac-



Fig. 7. Section in the structure through the (010) plane.

curacy but it does not obey definite laws. The N(3) atom motion is the least anisotropic and the least amplified. Amplitudes increase for the atoms at the bottom of an ethyl group.

No attempt to correct bond lengths for thermal motion has succeeded. The approximations of riding motion and of independant vibrations have been used but it may necessitate having to account for an important libration.

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# The Crystal Structure of In<sub>6</sub>Se<sub>7</sub>

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The crystal structure of  $In_6Se_7$  has been determined from powder and single-crystal data and is shown to be isomorphous with that of  $In_6S_7$ . The unit cell, which contains two formula units, is monoclinic with a=9.430, b=4.063, c=18.378 Å,  $\beta=109.34^\circ$  and the space group is  $P2_1$ . Basically the structure consists of two separate sections of almost cubic close-packed arrays of Se atoms with In atoms in octahedral coordination, the two sections having equivalent directions at  $61^\circ$  to each other.

#### Introduction

In their investigations of the phases existing in the In-Se system, Slavnova, Luzhnaya & Medvedeva (1963) and Slavnova & Eliseev (1963) described a black crystalline phase to which they ascribed the formula  $In_5Se_6$ .

This phase has been examined by the author as part of an X-ray crystallographic examination of the In–Se system. The structural analysis of the phase shows that